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(Article begins on next page)

THE TETRAHEDRAL MOTIF IN METAL STRUCTURES

David Turnbull and Michael J. Aziz
Division of Engineering and Applied Sciences
Harvard University
Cambridge, MA 02138

Abstract

F.C. Frank's demonstration of the stability of the icosahedral configuration to either of the close packed ones in thirteen-atom clusters of mutually attracting hard spheres renewed interest in the occurrence of tetrahedral configurations in metal alloy structures. This paper reviews some of these developments.

Consider the packing of mutually attracting, by central forces of hard spheres. In two dimensions the densest, and presumably lowest energy, packing unit is an equilateral triangular configuration which, when replicated, can fill space to form a hexagonal structure.

In contrast, in three dimensions, the configuration with lowest energy and highest density is a regular tetrahedron which, by itself, cannot be replicated to fill space and form a close-packed crystalline structure. The problem then arises of how these tetrahedral configurations, energetically preferred at short range, are incorporated in, or affect, three-dimensional structures. In this paper, various ways of including tetrahedral configurations into liquid, glass, and crystalline structures will be reviewed. Some excellent earlier treatments of this topic are:

M. Hoare [1]

F.C. Frank and J. Kasper [2]

F. Yonazawa [3]

J.F. Sadoc and R. Mosseri [4]

L. Bartell [5]

D.R. Nelson and F. Spaepen [6]

Hume-Rothery-Anderson model would predict. The melting temperature of the ordered intermetallic phase with composition $\sim\text{Au}_4\text{Si}$ is only 5°K below that of the actual eutectic.

Deposition from the pure metal vapor or dilute solution onto cold substrates might, initially, result in the formation of icosahedral, or other polytetrahedral, clusters (these have been labeled "amorphons"). These amorphons would continue to enlarge, and upon mutual impingement form a microquasicrystalline or, perhaps, a glassy structure. There is the intriguing possibility that below a certain thickness a quasicrystalline structure might be preferred, thermodynamically, to a microcrystalline one at the same atomic density. However, we know of no experimental evidence of a direct transformation from a micro-crystalline to a quasi-crystalline state.

A considerable number of metal alloys form structures in which icosahedral groups of atoms, as defined by Dirichlet constructions, are prominent components of the overall structure. Kasper also noted other coordination polytetrahedra which often appear in these structures. Frank and Kasper then examined these phases and outlined the general principle for inclusion of tetrahedral configurations in alloy structures. These alloys have come to be known as "Frank-Kasper phases".

Since icosahedral configurations appear to be important components of liquid structure, it seems that their resistance to nucleation of crystals with major polytetrahedral groups might be much lower than that to simple close-packed structures. This comparison can be made roughly in terms of the scaled, with thermodynamic melting temperature, T_m , undercooling at the onset of measurable crystal nucleation, $(T_m - T_N)/T_m = a$. For most pure metals, a is of order 0.2 to 0.25. At this time, information on those metals in which polytetrahedral groupings are prominent in their structures is rather sparse. However, recent studies at the Institut für Raumsimulation in Cologne by D. Holland-Moritz, D.M. Herlach and their coworkers, have indicated that a , for phases with major icosahedral or other polytetrahedral configurations in their structures, do indeed exhibit a values (~ 0.10) substantially below those with crystalline close-packed structures. The Cologne results were obtained by observation of droplets which were levitated electromagnetically. Earlier observations by Bendersky and Ridder [9] on the formation

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F.C. Frank's
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Hume-Rothery-Anderson model would predict. The melting temperature of the ordered intermetallic phase with composition $\sim\text{Au}_4\text{Si}$ is only 5°K below that of the actual eutectic.

Deposition from the pure metal vapor or dilute solution onto cold substrates might, initially, result in the formation of icosahedral, or other polytetrahedral, clusters (these have been labeled "amorphons"). These amorphons would continue to enlarge, and upon mutual impingement form a microquasicrystalline or, perhaps, a glassy structure. There is the intriguing possibility that below a certain thickness a quasicrystalline structure might be preferred, thermodynamically, to a microcrystalline one at the same atomic density. However, we know of no experimental evidence of a direct transformation from a micro-crystalline to a quasi-crystalline state.

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of quasicrystals from suspended droplets of Al-Mn alloy indicated in the condensed structure an extremely high concentration ($\sim 10^{18}/\text{cm}^3$) of microquasicrystallites.

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References

- [1] Hoare, M.R., *J. Non-Crystalline Solids*, **31**, 157 (1978).
- [2] Frank, F.C. and J. Kasper, *Acta Cryst.*, **11**, 184 (1958); **12**, 483 (1959).
- [3] Yonazawa, F., *Solid State Physics*, **45**, 179 (1991).
- [4] Sadoc, J.F. and R. Mosseri, *Geometrical Frustration* (Cambridge University Press, 1999).
- [5] Bartell, L., *Ann. Rev. Phys. Chem.* **49**, 43 (1998).
- [6] Nelson, D.R. and F. Spaepen, *Solid State Physics*, **42**, 1 (1989).
- [7] Hoare, M.R. and P. Pal, *Adv. Phys.*, **20**, 161 (1971).
- [8] Hume-Rothery, W. and E. Anderson, *Phil. Mag.*, **5**, 383 (1960).
- [9] Bendersky, L.A. and S.D. Ridder, *J. Mater. Res.*, **1**, 405 (1986).

PHYSICS AND METALLURGY

W.L. Johnson, 138-78 Keck

The basic factors governing the development of complex multicomponent systems are surveyed. The key kinetic and thermodynamic factors in these systems are outlined. Finally, the formation of liquid/crystal mixtures to produce unique properties is described.

Glass Formation in Metals

The early work of Turnbull demonstrated that a thermodynamic local fluctuation in the liquid in the form of a critical nucleus depends on the height of nucleation barrier (the height of nucleation barrier) and a time scale for rearrangement in the molten alloy.

$$(1) \quad R = \text{nucleation rate/unit volume}$$

where $v \sim \eta^{-1}$ is the rearrangement coefficient in the melt and is a function of undercooling. In the Turnbull approximation:

$$(2) \quad W \sim (16\pi/3)\gamma^3/\Delta g^2 \sim$$

where Δs_f is the specific entropy of crystallization (both per unit volume and per unit area), and T_m the melting point (approximation).

The rate of atomic rearrangement in the melt and is a function of undercooling through the Stokes-Einstein relation. The early 20th century, has shown the temperature dependence: